# Notes on Hodge theory 

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## 1 Hodge theory for real manifolds

### 1.1 Hodge *-operator

Let $(M, g)$ be a $n$-dimensional closed oriented Riemann manifold with volume form $\Omega$. Locally, we choose an orthonormal frame $\omega^{1}, \cdots, \omega^{n}$ with respect to $g$ for the cotangent bundle and thus we can write $\Omega=\omega^{1} \wedge \cdots \wedge \omega^{n}$. Denote the space of global smooth $k$-forms as $\Omega^{k}(M):=\Gamma\left(M, \wedge^{k} T M\right)$.

For $\omega, \eta \in \Omega^{k}(M)$, we write locally

$$
\begin{aligned}
& \omega=\sum_{i_{1}<\cdots<i_{k}} \omega_{i_{1}, \cdots, i_{k}} \omega^{i_{1}} \wedge \cdots \wedge \omega^{i_{k}} \\
& \eta=\sum_{i_{1}<\cdots<i_{k}} \eta_{i_{1}, \cdots, i_{k}} \omega^{i_{1}} \wedge \cdots \wedge \omega^{i_{k}} .
\end{aligned}
$$

Pointwisely, we define

$$
\langle\omega, \eta\rangle:=\sum_{i_{1}<\cdots<i_{k}} \omega_{i_{1}, \cdots, i_{k}} \eta_{i_{1}, \cdots, i_{k}} .
$$

It is easy to check $\langle\omega, \eta\rangle$ is a globally defined smooth function on $M$ and it is independent of the choice of orthonormal frames. So it is reasonable to define

$$
(\omega, \eta):=\int_{M}\langle\omega, \eta\rangle \Omega
$$

This gives an inner product on $\Omega^{*}(M)=\oplus_{k=0}^{n} \Omega^{k}(M)$. Note that $(-,-)$ is positive definite because $\Omega$ is nowhere vanishing.

Hodge $*$-operator $*: \Omega^{k}(M) \rightarrow \Omega^{n-k}(M)$ is defined by finding the 'complement' of the $k$-form in $\Omega$. More precisely, if

$$
\omega=\sum_{i_{1}<\cdots<i_{k}} \omega_{i_{1}, \cdots, i_{k}} \omega^{i_{1}} \wedge \cdots \wedge \omega^{i_{k}},
$$

then

$$
* \omega=\sum_{i_{1}<\cdots<i_{k}} \varepsilon_{i_{1}, \cdots, i_{k}} \omega_{i_{1}, \cdots, i_{k}} \omega^{1} \wedge \cdots \wedge \hat{\omega^{i_{1}}} \wedge \cdots \wedge \hat{\omega^{i_{k}}} \wedge \cdots \omega^{k},
$$

where

$$
\varepsilon_{i_{1}, \cdots, i_{k}}=(-1)^{i_{1}+\cdots+i_{k}+1+\cdots+k} .
$$

Proposition 1. Suppose $\omega, \eta \in \Omega^{k}(M)$. Then

1. $* 1=\Omega$,
2. $* \Omega=1$,
3. $* * \omega=(-1)^{k(n-k)} \omega$,
4. $\omega \wedge * \eta=\langle\omega, \eta\rangle \Omega$,
5. $\langle * \omega, * \eta\rangle=\langle\omega, \eta\rangle$.

Proof. All is straightforward from definition. We only prove 5 for example:

$$
\langle * \omega, * \eta\rangle \Omega=* \omega \wedge * * \eta=(-1)^{k(n-k)} * \omega \wedge \eta=\eta \wedge * \omega=\langle\eta, \omega\rangle \Omega .
$$

Then 5 follows from the fact: $\Omega$ is nowhere vanishing.
Recall $d: \Omega^{k-1}(M) \rightarrow \Omega^{k}(M)$. Define $\delta=(-1)^{n(k-1)+1} * d *: \Omega^{k}(M) \rightarrow \Omega^{k-1}(M)$.
Proposition 2. $\delta$ is the adjoint operator of $d$ with respect to $(-,-)$, i.e.,

$$
(d \omega, \eta)=(\omega, \delta \eta)
$$

for $\omega \in \Omega^{k-1}(M), \eta \in \Omega^{k}(M)$.
Proof. By direct computation,

$$
\begin{aligned}
d(\omega \wedge * \eta) & =d \omega \wedge * \eta+(-1)^{k-1} \omega \wedge d * \eta \\
& =d \omega \wedge * \eta+(-1)^{n(k-1)} \omega \wedge * * d * \eta \\
& =d \omega^{*} \eta-\omega \wedge * \delta \eta
\end{aligned}
$$

By Stoke's theorem,

$$
(d \omega, \eta)=\int_{M} d \omega^{*} \eta=\int_{M} \omega \wedge * \delta \eta=(\omega, \delta \eta)
$$

### 1.2 Harmonic forms and Hodge decomposition

Definition 1. $\Delta:=d \delta+\delta d: \Omega^{k}(M) \rightarrow \Omega^{k}(M)$ is called Hodge-Laplace operator. If $\omega \in \Omega^{k}(M)$ satisfies $\Delta \omega=0$, then $\omega$ is called a harmonic form.

Proposition 3. Hodge-Laplace operator satisfies:

1. $\Delta$ is self-adjoint, i.e., $(\Delta \omega, \eta)=(\omega, \Delta \eta)$ for all differtial forms.
2. $\Delta$ is positive, i.e., $(\Delta \omega, \omega) \geq 0$ and the equality holds if and only if $\Delta \omega=0$.
3. $* \Delta=\Delta *$.

Proof. To show $\Delta$ is self-adjoint, it suffices to assume $\omega$ and $\eta$ are both $k$-forms. Then

$$
\begin{aligned}
(\Delta \omega, \eta) & =(d \delta \omega, \eta)+(\delta d \omega, \eta) \\
& =(\delta \omega, \delta \eta)+(d \omega, d \eta) \\
& =(\omega, \Delta \eta) .
\end{aligned}
$$

Note that the above identity gives

$$
(\Delta \omega, \omega)=(\delta \omega, \delta \omega)+(d \omega, d \omega) \geq 0
$$

and $(\Delta \omega, \omega)=0$ if and only if $\delta \omega=0$ and $d \omega=0$, if and only if $\Delta \omega=0$.

To show $\Delta$ commutes with $*$, assume $\omega$ is a $k$-form, then

$$
* \delta \omega=(-1)^{n(k-1)+1} * * d * \omega=(-1)^{k} d * \omega .
$$

Similarly, $\delta * \omega=(-1)^{k+1} * d \omega$. So

$$
* d \delta \omega=(-1)^{k} \delta * \delta \omega=\delta d * \omega
$$

Similarly, $* \delta d=d \delta *$. Thus,

$$
* \Delta=* d \delta+* \delta d=\delta d *+d \delta *=\Delta *
$$

As we can see in the proof, $\omega$ is a harmonic form if and only if $d \omega=0$ and $\delta \omega=0$.
Denote $\mathcal{H}^{k}(M)$ as the vector space of harmonic $k$-forms on $M$. We will see $\mathcal{H}^{k}(M)$ is actually isomorphic to the de Rham cohomology of $M$.

Proposition 4. Suppose $\omega \in \mathcal{H}^{k}(M)$.

1. $\omega$ has minimal norm in the de Rham cohomology class $[\omega]$. More precisely, for any $(k-1)$-form $\eta$, $(\omega+d \eta, \omega+d \eta) \geq(\omega, \omega)$, and the equality holds if and only if $d \eta=0$.
2. $* \omega \in \mathcal{H}^{n-k}(M)$.

Proof. Suppose $\eta$ is a $(k-1)$-form. Then

$$
\begin{aligned}
(\omega+d \eta, \omega+d \eta) & =(\omega, \omega)+2(\omega, d \eta)+(d \eta, d \eta) \\
& =(\omega, \omega)+2(\delta \omega, \eta)+(d \eta, d \eta) \\
& =(\omega, \omega)+(d \eta, d \eta) \geq(\omega, \omega)
\end{aligned}
$$

The equality holds if and only if $d \eta=0 . * \omega \in \mathcal{H}^{n-k}(M)$ follows immediately from $\Delta *=* \Delta$.
Theorem 1 (Poincaré duality). Hodge *-operator gives an isomorphism $\mathcal{H}^{k}(M) \cong \mathcal{H}^{n-k}(M)$.
Proof. *: $\mathcal{H}^{k}(M) \rightarrow \mathcal{H}^{n-k}(M)$ is an isomorphism because $* *=(-1)^{k(n-k)}$.
To relate $\mathcal{H}^{k}(M)$ with the usual de Rham cohomology, we introduce our main result in the real case.

Theorem 2 (Hodge decomposition). There exists an isomorphism of vector spaces:

$$
\Omega^{k}(M)=\mathcal{H}^{k}(M) \oplus d \Omega^{k-1}(M) \oplus \delta \Omega^{k+1}(M)
$$

More precisely, for any $k$-form $\omega$, there exists a unique decomposition

$$
\omega=\omega_{h}+d \sigma+\delta \tau
$$

where $\omega_{h} \in \mathcal{H}^{k}(M), \sigma \in \Omega^{k-1}(M), \tau \in \Omega^{k+1}(M)$. When $\omega$ is closed, the decomposition reduces to

$$
\omega=\omega_{h}+d \sigma
$$

Theorem 3 (Poincaré duality for de Rham cohomology). We have an isomorphism $\mathcal{H}^{k}(M) \cong$ $H_{d R}^{k}(M)$, which gives Poincaré duality for de Rham cohomology

$$
H_{d R}^{k}(M) \cong H_{d R}^{n-k}(M)
$$

Proof. Define a linear map $\iota: \mathcal{H}^{k}(M) \rightarrow H_{d R}^{k}(M)$ by $\omega \mapsto[\omega] . \iota$ is injective by Proposition 4. $\iota$ is surjective by Hodge decomposition. So $\iota$ is an isomorphism.

## 2 Hodge theory for complex manifolds

### 2.1 Dolbeault cohomology

Let $X$ be a $n$-dimensional complex manifold. Denote $\Omega_{X}^{p}$ as the sheaf of holomorphic $p$-forms on $X$ and $\mathcal{A}_{X, \mathbb{C}}^{k}$ as the sheaf of complex $k$-forms on $X$. Recall the decomposition of sheaves

$$
\mathcal{A}_{X, \mathbb{C}}^{k}=\bigoplus_{p+q=k} \mathcal{A}_{X}^{p, q}
$$

and the differential

$$
\begin{aligned}
& \partial: \mathcal{A}_{X}^{p, q} \rightarrow \mathcal{A}_{X}^{p+1, q} \\
& \bar{\partial}: \mathcal{A}_{X}^{p, q} \rightarrow \mathcal{A}_{X}^{p, q+1},
\end{aligned}
$$

where $\mathcal{A}_{X}^{p, q}$ is the sheaf of forms of type $(p, q)$ on $X$. Recall these sheaves $\mathcal{A}^{p, q}$ are acyclic, i.e., have trivial higher cohomology from partition of unity. Then Dolbeault cohomology with respect to differential forms is defined by

$$
H^{p, q}(X):=H^{q}\left(\mathcal{A}^{p, \bullet}(X), \bar{\partial}\right)=\frac{\operatorname{Ker}\left(\bar{\partial}: \mathcal{A}^{p, q}(X) \rightarrow \mathcal{A}^{p, q+1}(X)\right)}{\operatorname{Im}\left(\bar{\partial}: \mathcal{A}^{p, q-1}(X) \rightarrow \mathcal{A}^{p, q}(X)\right)}
$$

In fact, the Dolbeault cohomology is isomorphic to sheaf cohomology of $\Omega_{X}^{p}$, i.e.,

$$
H^{p, q}(X) \cong H^{q}\left(X, \Omega_{X}^{p}\right)
$$

To see this, it suffices to see the acyclic resolution

$$
0 \rightarrow \Omega_{X}^{p} \rightarrow \mathcal{A}^{p, 0} \xrightarrow{\overline{\bar{c}}} \mathcal{A}^{p, 1} \rightarrow \cdots
$$

from $\bar{\partial}$-Poincaré lemma.
Let $E$ be a complex vector bundle over $X$ and $\mathcal{A}^{p, q}(E)$ denote the sheaf defined by

$$
U \mapsto \mathcal{A}^{p, q}(U, E):=\Gamma\left(U, \wedge^{p, q} X \otimes E\right)
$$

where the tensor product is taken over $\mathcal{O}_{X}$. Locally, a section $\alpha$ of $\mathcal{A}^{p, q}(E)$ can be written as $\alpha=\sigma \alpha_{i} \otimes s_{i}$ with $\alpha$ and $s_{i}$ local sections of $\mathcal{A}_{X}^{p, q}$ and $E$ respectively.

Lemma 1. Suppose $E$ is a holomorphic vector bundle. There exists a natural $\mathbb{C}$-linear operator $\bar{\partial}_{E}: \mathcal{A}^{p, q}(E) \rightarrow \mathcal{A}^{p, q+1}(E)$ with $\bar{\partial}_{E}^{2}=0$ and which satisfies the Leibniz rule

$$
\bar{\partial}_{E}(f \cdot \alpha)=\bar{\partial}(f) \wedge \alpha+f \bar{\partial}_{E}(\alpha)
$$

Proof. Choose a local trivialization $s=\left(s_{1}, \cdots, s_{r}\right)$ of $E$ and write $\alpha \in \mathcal{A}^{p, q}(E)$ locally as $\alpha=$ $\sum \alpha_{i} \otimes s_{i}$, where $\alpha \in \mathcal{A}_{X}^{p, q}$. Define

$$
\bar{\partial}_{E} \alpha:=\sum \bar{\partial}\left(\alpha_{i}\right) \otimes s_{i}
$$

Suppose we choose another holomorphic trivialization $s^{\prime}=\left(s_{1}^{\prime}, \cdots, s_{r}^{\prime}\right)$ and obtain an operator $\bar{\partial}_{E}^{\prime}$. Let $s_{i}=\sum_{j} \psi_{i j} s_{j}^{\prime}$, where $\psi_{i j}$ is the holomorphic transition function. Then

$$
\begin{aligned}
\bar{\partial}_{E}^{\prime} \alpha & =\bar{\partial}_{E}^{\prime}\left(\sum \alpha_{i} \otimes \sum_{j} \psi_{i j} s_{j}^{\prime}\right) \\
& =\sum_{i, j} \bar{\partial}\left(\alpha_{i} \psi_{i j}\right) \otimes s_{j}^{\prime} \\
& =\sum_{i, j} \bar{\partial}\left(\alpha_{i}\right) \psi_{i j} \otimes s_{j}^{\prime}=\bar{\partial}_{E}(\alpha) .
\end{aligned}
$$

So $\bar{\partial}_{E}=\bar{\partial}_{E}^{\prime}$ is independent of the choice of local trivialization. Therefore, $\bar{\partial}_{E}^{2}=0$ since $\bar{\partial}^{2}=0$. From Leibniz rule of $\bar{\partial}$,

$$
\begin{aligned}
\bar{\partial}_{E}(f \cdot \alpha) & =\bar{\partial}_{E}\left(\sum f \alpha_{i} \otimes s_{i}\right) \\
& =\sum \bar{\partial}\left(f \alpha_{i}\right) \otimes s_{i} \\
& =\sum\left(\bar{\partial}(f) \wedge \alpha_{i}+f \bar{\partial}\left(\alpha_{i}\right)\right) \otimes s_{i} \\
& =\bar{\partial}(f) \wedge \alpha+f \bar{\partial}_{E}(\alpha) .
\end{aligned}
$$

The above lemma gives a complex $\left(\mathcal{A}^{p, \bullet}(X, E), \bar{\partial}_{E}\right)$, whose cohomology is called Dolbeault cohomology of the holomorphic vector bundle $E$ :

$$
H^{p, q}(X, E):=H^{q}\left(\mathcal{A}^{p, \bullet}(X, E), \bar{\partial}_{E}\right)=\frac{\operatorname{Ker}\left(\bar{\partial}_{E}: \mathcal{A}^{p, q}(X, E) \rightarrow \mathcal{A}^{p, q+1}(X, E)\right)}{\operatorname{Im}\left(\bar{\partial}_{E}: \mathcal{A}^{p, q-1}(X, E) \rightarrow \mathcal{A}^{p, q}(X, E)\right)} .
$$

Similarly, Dolbeault cohomology of holomorphic vector bundles is isomorphic to sheaf cohomology:

$$
H^{p, q}(X, E) \cong H^{q}\left(X, E \otimes \Omega_{X}^{p}\right)
$$

which follows from acyclic resolution of $E \otimes \Omega_{X}^{p}$ :

$$
0 \rightarrow E \otimes \Omega_{X}^{p} \rightarrow \mathcal{A}^{p, 0}(E) \xrightarrow{\bar{\partial}_{E}} \mathcal{A}^{p, 1}(E) \rightarrow \cdots
$$

To summarize what we obtain:
Dolbeault cohomology of a holomorphic vector bundle $E=$ sheaf cohomology of $E \otimes \Omega_{X}^{p}$.

### 2.2 Hermitian and Kähler structure on complex manifolds

We briefly recall some definition in this subsection.
Let $X$ be a complex manifold with almost complex structure $I$ and complex dimension $n$. A Riemann metric $g$ on $X$ is an hermitian structure if for any point $x \in X$, the scalar product $g_{x}$ is compatible with $I$, i.e.,

$$
g_{x}(I v, I w)=g_{x}(v, w)
$$

The induced real $(1,1)$-form $\omega:=g(I(),())$ is called the fundamental form of hermitian manifold $(X, g)$. After complexification, the fundamental form $\omega$ is locally of the form

$$
\omega=\frac{i}{2} \sum_{i, j=1}^{n} h_{i j} d z_{i} \wedge d \overline{z_{j}},
$$

where $\left(h_{i j}(x)\right)$ is a positive definition matrix for each $x \in X$. It is not difficult to see that the hermitian structure is uniquely determined by $I$ and $\omega$. The hermitian structure $g$ is called a Kähler structure if $\omega$ is closed. Denote the hermitian extension of $g$ by $g_{\mathbb{C}}$.

The Hodge $*$-operator $*: \bigwedge_{\mathbb{C}}^{k} X \rightarrow \bigwedge_{\mathbb{C}}^{2 n-k} X$ is similar to the one defined for Riemann manifolds $(X, g)$ with the natural volume form $\Omega$. More precisely, $*$ is defined by $\alpha \wedge * \bar{\beta}=g_{\mathbb{C}}(\alpha, \beta) \Omega$. When restricted to $\bigwedge_{\mathbb{C}}^{k} X, *$ reduces to the usual Hodge $*$-operator for Riemann manifolds.

With Hodge $*$-operator, we can define several adjoint operators. Regard $(X, g)$ as a real Riemann manifold with natural volume form $\Omega$. Since $X$ has even dimension, the adjoint operator $d^{*}=\delta$ is exactly $d^{*}=-* \circ d \circ *$. Analogously, one defines $\partial^{*}$ and $\bar{\partial}^{*}$ as $\partial^{*}:=-* \circ \bar{\partial} \circ *$ and $\bar{\partial}^{*}=-* \circ \partial \circ *$ to make $d^{*}=\partial^{*}+\bar{\partial}^{*}$ valid.

Therefore, it is natural to define the Laplacian operator associated to $d, \partial, \bar{\partial}: \Delta:=d^{*} d+d d^{*}$, $\Delta_{\partial}=\partial^{*} \partial+\partial \partial^{*}, \Delta_{\bar{\partial}}=\bar{\partial}^{*} \bar{\partial}+\bar{\partial} \bar{\partial}^{*}$. Note that $*: \mathcal{A}^{p, q}(X) \rightarrow \mathcal{A}^{n-q, n-p}(X), \partial^{*}: \mathcal{A}^{p, q}(X) \rightarrow$ $\mathcal{A}^{p-1, q}(X), \bar{\partial}^{*}: \mathcal{A}^{p, q}(X) \rightarrow \mathcal{A}^{p, q-1}(X)$ and $\Delta_{\partial}, \Delta_{\bar{\partial}}$ perserve bidegrees.

### 2.3 Hodge decomposition for Kähler manifolds

Syppose $X$ is a complex manifold with an hermitian structure $g$ and natural fundamental form $\Omega$. The hermitian extension of $g$ is denoted by $g_{\mathbb{C}}$. Define an hermitian product on $\mathcal{A}_{\mathbb{C}}^{*}(X)$ by

$$
(\alpha, \beta):=\int_{X} g_{\mathbb{C}}(\alpha, \beta) * 1=\int_{X} g_{\mathbb{C}}(\alpha, \beta) \Omega .
$$

Note that the value of $g$ on $\mathcal{A}^{*}(X)$ is the exactly the inner product defined in the first section.
With respect to the hermitian product $(-,-)$. The degree decomposition

$$
\mathcal{A}_{\mathbb{C}}^{*}(X)=\bigoplus_{k} \mathcal{A}_{\mathbb{C}}^{k}(X)
$$

and the bidegree decomposition

$$
\mathcal{A}_{\mathbb{C}}^{k}(X)=\bigoplus_{p+q=k} \mathcal{A}^{p, q}(X)
$$

are both orthogonal decompositions and each component $\mathcal{A}^{p, q}(X)$ is an infinite dimensional normed vector space with scalar product $(-,-)$ and the induced norm $\|\alpha\|^{2}=(\alpha, \alpha)$.

Proposition 5. Suppose $X$ is a closed hermitian manifold. Then with respect to $(-,-)$, the operators $d^{*}, \partial^{*}, \bar{\partial}^{*}$ are actually adjoint operators of $d, \partial, \bar{\partial}$.
Proof. We give a proof for $\partial$ for example. For $\alpha \in \mathcal{A}^{p-1, q}(X)$ and $\beta \in \mathcal{A}^{p, q}(X)$,

$$
\begin{aligned}
(\partial \alpha, \beta) & =\int_{X} g_{\mathbb{C}}(\partial \alpha, \beta) * 1=\int_{X} \partial \alpha \wedge * \bar{\beta} \\
& =\int_{X} \partial(\alpha \wedge * \bar{\beta})-(-1)^{p+q-1} \int_{X} \alpha \wedge \partial(* \bar{\beta})
\end{aligned}
$$

Since $\alpha \wedge * \bar{\beta}$ is of bidegree $(n-1, n), \partial(\alpha \wedge * \bar{\beta})=d(\alpha \wedge * \bar{\beta})$. So

$$
\int_{X} \partial(\alpha \wedge * \bar{\beta})=0
$$

by Stokes' theorem. Using $*^{2}=(-1)^{k}$ on $\mathcal{A}^{k}(X)$, we compute

$$
\int_{X} \alpha \wedge \partial(* \bar{\beta})=(-1)^{2 n-(p+q)+1} \int_{X} g_{\mathbb{C}}\left(\alpha,-\partial^{*} \beta\right) * 1=(-1)^{2 n-(p+q)}\left(\alpha, \partial^{*} \beta\right) .
$$

So $(\partial \alpha, \beta)=\left(\alpha, \partial^{*} \beta\right)$.
In the case of real manifolds, we interpret de Rham cohomology by harmonic forms. In the complex case, we will apply similar approach. For the differential $d$, the spaces of harmonic $k$ forms and ( $p, q$ )-forms (which are defined similarly in the real case) are denoted by $\mathcal{H}^{k}(X, g)$ and $\mathcal{H}^{p, q}(X, g)$. For $\partial$ and $\bar{\partial}$, we have analogous definition.

Definition 2. A $k$-form is called $\bar{\partial}$-harmonic if $\Delta_{\bar{\partial}} \alpha=0$ and define the spaces of $\bar{\partial}$-harmonic $k$-forms and $(p, q)$-forms by $\mathcal{H}_{\bar{\partial}}^{k}(X, g)$ and $\mathcal{H}_{\bar{\partial}}^{p, q}(X, g)$. $\partial$-harmonic forms are analogous.

Proposition 6. Suppose $(X, g)$ is a closed hermitian manifold. A form $\alpha$ is $\bar{\partial}$-harmonic (resp. $\partial$-harmonic) if and only if $\bar{\partial} \alpha=\bar{\partial}^{*} \alpha=0$ (resp. $\partial \alpha=\partial^{*} \alpha=0$ ).

Proof. The $\bar{\partial}$ case follows from the identity

$$
\begin{aligned}
\left(\Delta_{\bar{\partial}} \alpha, \alpha\right) & =\left(\bar{\partial}^{*} \bar{\partial} \alpha+\bar{\partial} \bar{\partial}^{*} \alpha, \alpha\right) \\
& =\left\|\bar{\partial}^{*}(\alpha)\right\|^{2}+\|\bar{\partial}\|^{2}
\end{aligned}
$$

Proposition 7. Suppose $(X, g)$ is an hermitian manifold. Then

1. $\mathcal{H}_{\bar{\partial}}^{k}(X, g)=\bigoplus_{p+q=k} \mathcal{H}_{\bar{\partial}}^{p, q}(X, g)$ and $\mathcal{H}_{\partial}^{k}(X, g)=\bigoplus_{p+q=k} \mathcal{H}_{\partial}^{p, q}(X, g)$.
2. If $(X, g)$ is Kähler, then both decompositions coincide with $\mathcal{H}^{k}(X, g)_{\mathbb{C}}=\bigoplus_{p+q=k} \mathcal{H}^{p, q}(X, g)$. In particular, $\mathcal{H}^{k}(X, g)_{\mathbb{C}}=\mathcal{H}_{\bar{\partial}}^{k}(X, g)=\mathcal{H}_{\partial}^{k}(X, g)$.

Proof. Suppose $\alpha=\sum \alpha^{p, q}$ is the bidegree decomposition of a $\bar{\partial}$-harmonic form $\alpha$. Then

$$
0=\sum \Delta_{\bar{\partial}}\left(\alpha^{p, q}\right)
$$

is also a bidegree decomposition, which implies $\Delta_{\bar{\partial}}\left(\alpha^{p, q}\right)=0$ for all $p, q$. The proof of $\partial$ decomposition is analogous.

The second assertion follows from the identity $\Delta_{\partial}=\Delta_{\bar{\partial}}=\frac{1}{2} \Delta$ on a Kähler manifold.
Theorem 4 (Serre duality for harmonic forms). Suppose $(X, g)$ is a compact connected hermitian manifold. Then the pairing

$$
\mathcal{H}_{\bar{\partial}}^{p, q}(X, g) \times \mathcal{H}_{\bar{\partial}}^{n-p, n-q}(X, g) \rightarrow \mathbb{C}, \quad(\alpha, \beta) \mapsto \int_{X} \alpha \wedge \beta
$$

is non-degenerate. This yields an isomorphism

$$
\mathcal{H}_{\bar{\partial}}^{p, q}(X, g) \cong \mathcal{H}_{\bar{\partial}}^{n-p, n-q}(X, g)^{*}
$$

Proof. Suppose $0 \neq \alpha \in \mathcal{H}_{\bar{\partial}}^{p, q}(X, g)$. Then

$$
\int_{X} \alpha \wedge * \bar{\alpha}=\|\alpha\|^{2}>0
$$

implies the pairing is non-degenerate.
Our main result is:
Theorem 5 (Hodge decomposition of harmonic forms). Let $(X, g)$ be a compact hermitian manifold. Then there exists two natural orthogonal decompositions

$$
\begin{aligned}
\mathcal{A}^{p, q}(X) & =\partial \mathcal{A}^{p-1, q}(X) \oplus \mathcal{H}_{\partial}^{p, q}(X, g) \oplus \partial^{*} \mathcal{A}^{p+1, q}(X) \\
\mathcal{A}^{p, q}(X) & =\bar{\partial} \mathcal{A}^{p-1, q}(X) \oplus \mathcal{H}_{\bar{\partial}}^{p, q}(X, g) \oplus \bar{\partial}^{*} \mathcal{A}^{p+1, q}(X)
\end{aligned}
$$

Moreover, $\mathcal{H}^{p, q}(X, g)_{\mathbb{C}}$ are all finite dimensional and if $X$ is a Kähler manifold, then $\mathcal{H}^{p, q}(X, g)_{\mathbb{C}}=$ $\mathcal{H}_{\partial}^{p, q}(X, g)=\mathcal{H}_{\bar{\partial}}^{p, q}(X, g)$.

The most nontrivial part is the existence of decomposition. We will not prove it in our note. The significance of Kähler condition is that we can forget all about ' $\bar{\partial}$ or $\bar{\partial}$-harmonic', and replace them by 'harmonic'.

Corollary 1. Suppose $(X, g)$ is a compact hermitian manifold. Then the canonical map $\mathcal{H}_{\bar{\rho}}^{p, q}(X, g) \rightarrow$ $H^{p, q}(X)$ is an isomorphism.

Proof. The canonical map $\mathcal{H}_{\bar{\partial}}^{p, q}(X, g) \rightarrow H^{p, q}(X)$ is given by $\alpha \mapsto[\alpha]$, where $[\alpha]$ is a cohomology class of $\alpha$. This map is injective because any harmonic coboundary $\bar{\partial} \beta$ must satisfy $\bar{\partial} * \bar{\partial} \beta=0$. But $0=\left(\bar{\partial}^{*} \bar{\partial} \beta, \beta\right)=(\bar{\partial} \beta, \bar{\partial} \beta)=\|\bar{\partial} \beta\|^{2}=0$ implies $\bar{\partial} \beta=0$. To show this map is surjective, it suffices to show any $\bar{\partial}$-closed $(p, q)$-form $\beta$ must be cohomological to a $\bar{\partial}$-harmonic form. By Hodge decomposition, write $\beta=\bar{\partial} \beta_{1}+\beta_{h}+\bar{\partial}^{*} \beta_{2}$, where $\beta_{h}$ is harmonic. Since $\bar{\partial} \beta=0$, it follows $\bar{\partial} \bar{\partial}^{*} \beta_{2}=0$, which implies $\bar{\partial}^{*} \beta_{2}$ by a similar argument in the injective case.

The proof shows the Hodge decomposition for closed forms does not contain the terms $\partial^{*} \mathcal{A}^{p+1, q}(X)$ or $\bar{\partial}^{*} \mathcal{A}^{p+1, q}(X)$, just as the case in real manifolds.

Corollary 2. Let $(X, g)$ be a compact Kähler manifold. Then there exists a decomposition

$$
H^{k}(X, \mathbb{C})=\bigoplus_{p+q=k} H^{p, q}(X)
$$

which does not depend on the Kähler structure.
Proof. Since $X$ is Kähler,

$$
H^{k}(X, \mathbb{C})=\mathcal{H}^{k}(X, g)_{\mathbb{C}}=\bigoplus_{p+q=k} \mathcal{H}^{p, q}(X, g)=\bigoplus_{p+q=k} H^{p, q}(X)
$$

To show the decomposition is independent of the Kähler structure, we choose two Kähler metric $g, g^{\prime}$ and $\alpha \in \mathcal{H}^{p, q}(X, g), \alpha^{\prime} \in \mathcal{H}^{p, q}\left(X, g^{\prime}\right)$, which induce same elements in $H^{p, q}(X)$. So $\alpha$ and $\alpha^{\prime}$ differ by some $\bar{\partial} \gamma$, i.e., $\alpha^{\prime}=\alpha+\bar{\partial} \gamma$. Then $d \bar{\partial} \gamma=0$. By Hodge decomposition for $d$,

$$
\bar{\partial} \gamma=d \beta+\beta_{h}
$$

But $0=\left(\gamma, \bar{\partial}^{*} \beta_{h}\right)=\left(\bar{\partial} \gamma, \beta_{h}\right)=\left(\beta_{h}, \beta_{h}\right)$ implies $\beta_{h}=0$. So $\bar{\partial} \gamma \in d\left(\mathcal{A}_{\mathbb{C}}^{k-1}(X)\right)$ and $\alpha, \alpha^{\prime}$ induces the same de Rham cohomology class in $H^{k}(X, \mathbb{C})$.

## 3 Serre duality

In this section, we will give a generalized version of Hodge decomposition and Serre duality on a holomorphic vector bundle. Serre duality, together with Riemann-Roch theorem and Kodaira vanishing theorem, is significant in controlling the cohomology of holomorphic vector bundles. Most parts of this section are routine.

### 3.1 Hermitian structure on vector bundles

Suppose $M$ is a real manifold and $E \rightarrow M$ is a complex vector bundle.
Definition 3. An hermitian structure $h$ on $E \rightarrow M$ is an hermitian scalar product $h_{x}$ on each fiber $E(x)$ which depends differtiably on $x$, The pair $(E, h)$ is called an hermitian vector bundle.

If $\psi:\left.E\right|_{U} \cong U \times \mathbb{C}^{r}$ is a trivialization over some open subset $U$, then $h_{x}$ is given by a positivedefinite hermitian matrix $\left(h_{i j}(x)\right)$ for each $x \in U$. The definition says $\left(h_{i j}(x)\right)$ relies differtiably on $x \in U$.

Example 1. If $(X, g)$ is an hermitian manifold, then the tangent, cotangent and form bundles have natural hermitian structures. Moreover, if $(E, h)$ is an hermitian vector bundle over $(X, g)$, then $\bigwedge^{p, q} X \otimes E$ have natural hermitian structures.

Proposition 8. Every complex vector bundle admits an hermitian structure.
Proof. Choose an open covering $X=\bigcup U_{i}$ trivializing $E$ and glue the constant hermitian structure on the trivial bundles $U_{i} \times \mathbb{C}^{r}$ by means of partition of unity. The resulting product is hermitian because positive linear combination of hermitian product is again hermitian.

Now let $(X, g)$ be an hermitian manifold of complex dimension $n$ and $(E, h)$ is an hermitian vector bundle. Denote the induced hermitian structure on $\Lambda^{p, q} X \otimes E$ by $(-,-)$. We may interpret $h$ as a $\mathbb{C}$-antilinear isomorphism $h: E \cong E^{*}$.
Definition 4. Hodge $*$-operator $\bar{F}_{E}: \bigwedge^{p, q} X \otimes E \rightarrow \bigwedge^{n-p, n-q} X \otimes E^{*}$ is defined by

$$
\bar{*}_{E}(\varphi \otimes s)=\bar{*}_{E}(\varphi) \otimes h(s)=\overline{*(\varphi)} \otimes h(s)=*(\bar{\varphi}) \otimes h(s) .
$$

Hodge $*$-operator $\bar{*}_{E}$ is a $\mathbb{C}$-antilinear isomorphism that depends on $g$ and $h$. Similarly, we check easily

$$
(\alpha, \beta) * 1=\alpha \wedge \bar{*}_{E}(\beta)
$$

where $\wedge$ means taking usual wedge products in the form part and evaluation map in the bundle part. Moreover, $\bar{*}_{E} \circ \bar{*}_{E}=(-1)^{p+q}$ on $\bigwedge^{p, q} X \otimes E$. From now on, denote $\mathcal{A}^{p, q}(E)$ by sheaf of sections of $\bigwedge^{p, q} X \otimes E$. We will not distinguish $\mathcal{A}^{p, q}(E)$ and $\bigwedge^{p, q} X \otimes E$ from now on.

Define $\bar{\partial}_{E}: \mathcal{A}^{p, q}(E) \rightarrow \mathcal{A}^{p, q+1}(E)$ by $\bar{\partial}_{E}(\alpha \otimes s)=\bar{\partial}(\alpha) \otimes s$ and its adjoint operator $\bar{\partial}_{E}^{*}$ : $\mathcal{A}^{p, q}(E) \rightarrow \mathcal{A}^{p, q-1}(E)$ by $\bar{\partial}_{E}^{*}=-\bar{*}_{E^{*}} \circ \bar{\partial}_{E^{*}} \circ \bar{*}_{E}$. The Laplacian operator $\Delta_{E}: \mathcal{A}^{p, q}(E) \rightarrow \mathcal{A}^{p, q}(E)$ is defined by $\Delta_{E}=\bar{\partial}_{E}^{*} \bar{\partial}_{E}+\bar{\partial}_{E} \bar{\partial}_{E}^{*}$.

Definition 5. A section $\alpha \in \mathcal{A}^{p, q}(E)$ is called harmonic if $\Delta_{E}(\alpha)=0$. The space of harmonic forms is denoted by $\mathcal{H}^{p, q}(X, E)$.

Since $\bar{\star}_{E}$ commutes with $\Delta_{E}, \bar{\star}_{E}$ restricts to a $\mathbb{C}$-antilinear isomorphism $\bar{\star}_{E}: \mathcal{H}^{p, q}(X, E) \rightarrow$ $\mathcal{H}^{p, q}\left(X, E^{*}\right)$.

From now on, we suppose $(X, g)$ is a compact hermitian manifold. Define a hermitian scalar product on $\mathcal{A}^{p, q}(X, E)$ by

$$
(\alpha, \beta)=\int_{X}(\alpha, \beta) * 1
$$

where $(-,-)$ inside the integral is the pointwise hermitian inner product on $\mathcal{A}^{p, q}(X, E)$.
Proposition 9. 1. $\bar{\partial}_{E}^{*}$ is actually the adjoint operator of $\bar{\partial}_{E}$ and $\Delta_{E}$ is self-adjoint with respect to $(-,-)$.
2. $\alpha \in \mathcal{A}^{p, q}(X, E)$ is harmonic if and only if $\bar{\partial}_{E}(\alpha)=\bar{\partial}_{E}^{*}(\alpha)=0$.

Proof. Analogous to Proposition 2 and 3.

### 3.2 Serre duality on vector bundles

To give Serre duality on vector bundles, we have to give a generalized version of Hodge decomposition for vector bundles.

Theorem 6 (Hodge decomposition for vector bundles). Suppose ( $X, g$ ) is a compact hermitian manifold and $(E, h)$ is an hermitian vector bundle. We have Hodge decomposition

$$
\mathcal{A}^{p, q}(X, E)=\bar{\partial}_{E} \mathcal{A}^{p, q-1}(X, E) \oplus \mathcal{H}^{p, q}(X, E) \oplus \bar{\partial}_{E}^{*} \mathcal{A}^{p, q+1}(X, E)
$$

and $\mathcal{H}^{p, q}(X, E)$ is finite dimensional for each $p$ and $q$.

If we choose $E=\mathcal{O}_{X}$, the above theorem reduces to the usual Hodge decomposition.
Corollary 3. The natural map $\mathcal{H}^{p, q}(X, E) \rightarrow H^{p, q}(X, E)$ is an isomorphism. In particular, Dolbeault cohomology $H^{p, q}(X, E) \cong H^{q}\left(X, E \otimes \Omega_{X}^{p}\right)$ is finite dimensional.
Proof. To show the map is injective, choose a harmonic coboundary $\bar{\partial}_{E} \beta$. So

$$
0=\left(\bar{\partial}_{E}^{*} \bar{\partial}_{E} \beta, \beta\right)=\left(\bar{\partial}_{E} \beta, \bar{\partial}_{E} \beta\right)=\left\|\bar{\partial}_{E} \beta\right\|^{2}
$$

implies $\bar{\partial}_{E} \beta=0$.
To show the map is surjective, choose a closed element $\alpha \in \mathcal{A}^{p, q}(X, E)$ and by Hodge decomposition,

$$
\alpha=\bar{\partial}_{E} \alpha_{1}+\alpha_{h}+\bar{\partial}_{E}^{*} \alpha_{2} .
$$

Since $\bar{\partial}_{E} \alpha=0, \bar{\partial}_{E} \bar{\partial}_{E}^{*} \alpha_{2}=0$ and then $\bar{\partial}_{E}^{*} \alpha_{2}=0$. So $\alpha=\bar{\partial}_{E} \alpha_{1}+\alpha_{h}$, i.e., the image of $\alpha_{h} \in$ $\mathcal{H}^{p, q}(X, E)$ under the natural map is the cohomology class $[\alpha]$.

By this corollary, we see any (nonzero) Dolbeault cohomology class can be represented by a (nonzero) harmonic element.

Theorem 7 (Serre duality for vector bundles). Suppose $X$ is a compact complex manifold. For any holomorphic vector bundle $E$ on $X$, define a pairing

$$
H^{p, q}(X, E) \times H^{n-p, n-q}\left(X, E^{*}\right) \rightarrow \mathbb{C}
$$

by

$$
(\alpha, \beta) \mapsto \int_{X} \alpha \wedge \beta
$$

Then the pairing is non-degenerate.
Proof. This pairing is well-defined by Stokes' theorem. Choose hermitian structures $h$ and $g$ on $E$ and $X$. For any nonzero cohomology class in $H^{p, q}(X, E)$, we can find a nonzero harmonic element $\alpha \in \mathcal{H}^{p, q}(X, E)$ representing the given class. Define $\beta=\bar{\star}_{E} \alpha \in \mathcal{H}^{n-p, n-q}\left(X, E^{*}\right)$. We check

$$
\int_{X} \alpha \wedge \beta=\int_{X} \alpha \wedge \bar{*}_{E} \alpha=\int_{X}(\alpha, \alpha) * 1=\|\alpha\|^{2} \neq 0 .
$$

So this pairing is non-degenerate.
Corollary 4. For any holomorphic vector bundle $E$ over a compact complex manifold $X$, there exist $\mathbb{C}$-linear isomorphisms:

$$
H^{q}\left(X, E \otimes \Omega^{p}\right) \cong H^{p, q}(X, E) \cong H^{n-p, n-q}\left(X, E^{*}\right)^{*} \cong H^{n-q}\left(X, E^{*} \otimes \Omega^{n-p}\right)^{*}
$$

